

Modeling Photo-induced dynamical processes in massive parallel architectures

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OUTLINE

Introduction: framing the problem

“theoretical synchrotron: ETSF”

“Modeling new materials properties”

Theoretical background:

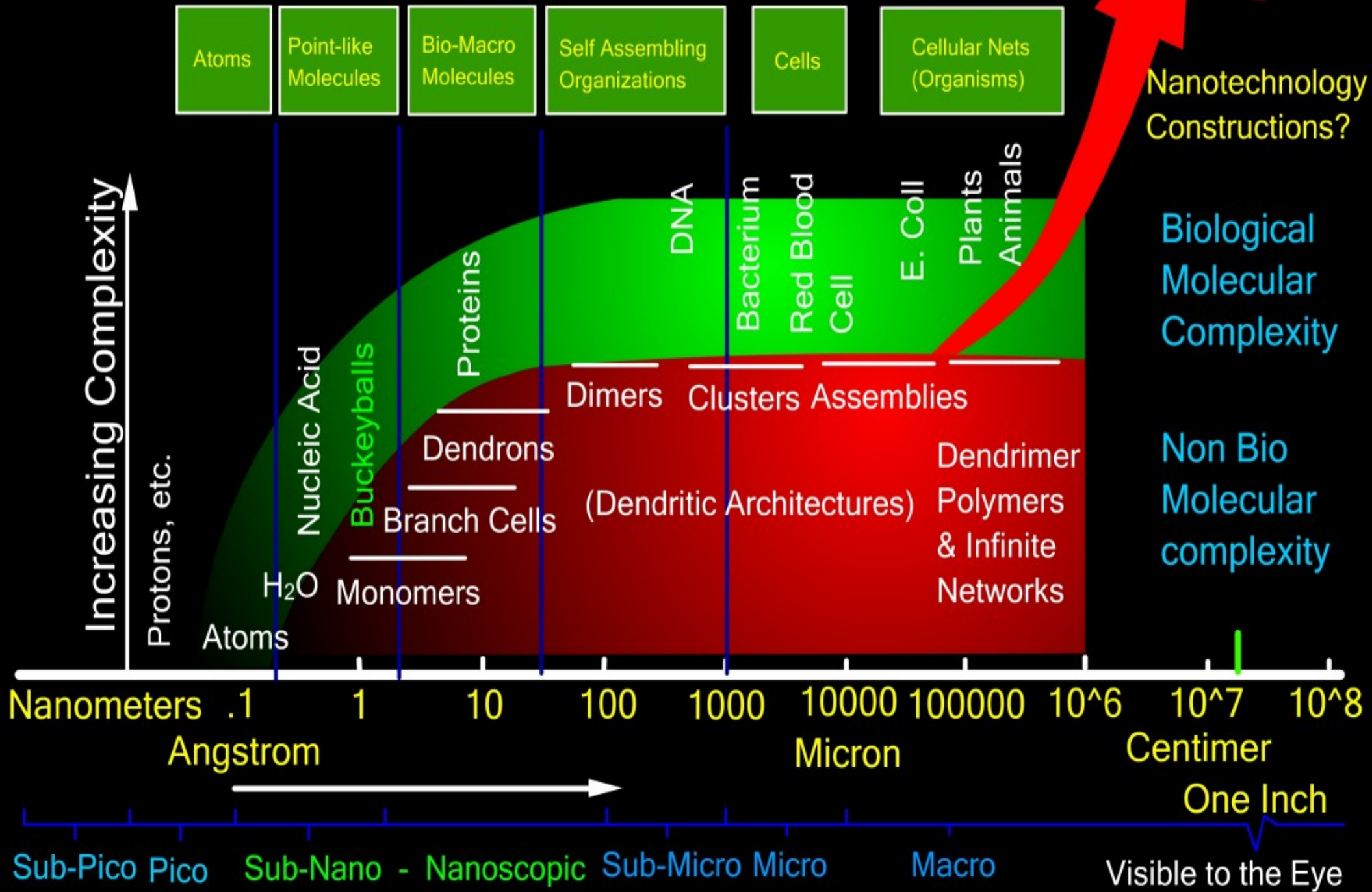
Some Applications:

- * (excite state dynamics and energy store/harvesting)
 - * photo-excitation light-harvesting carotenoid-porphyrin-C60
 - * **Chlorophylls: photosynthesis**
-
- * Optimal control of bond-breaking and
high harmonic generation

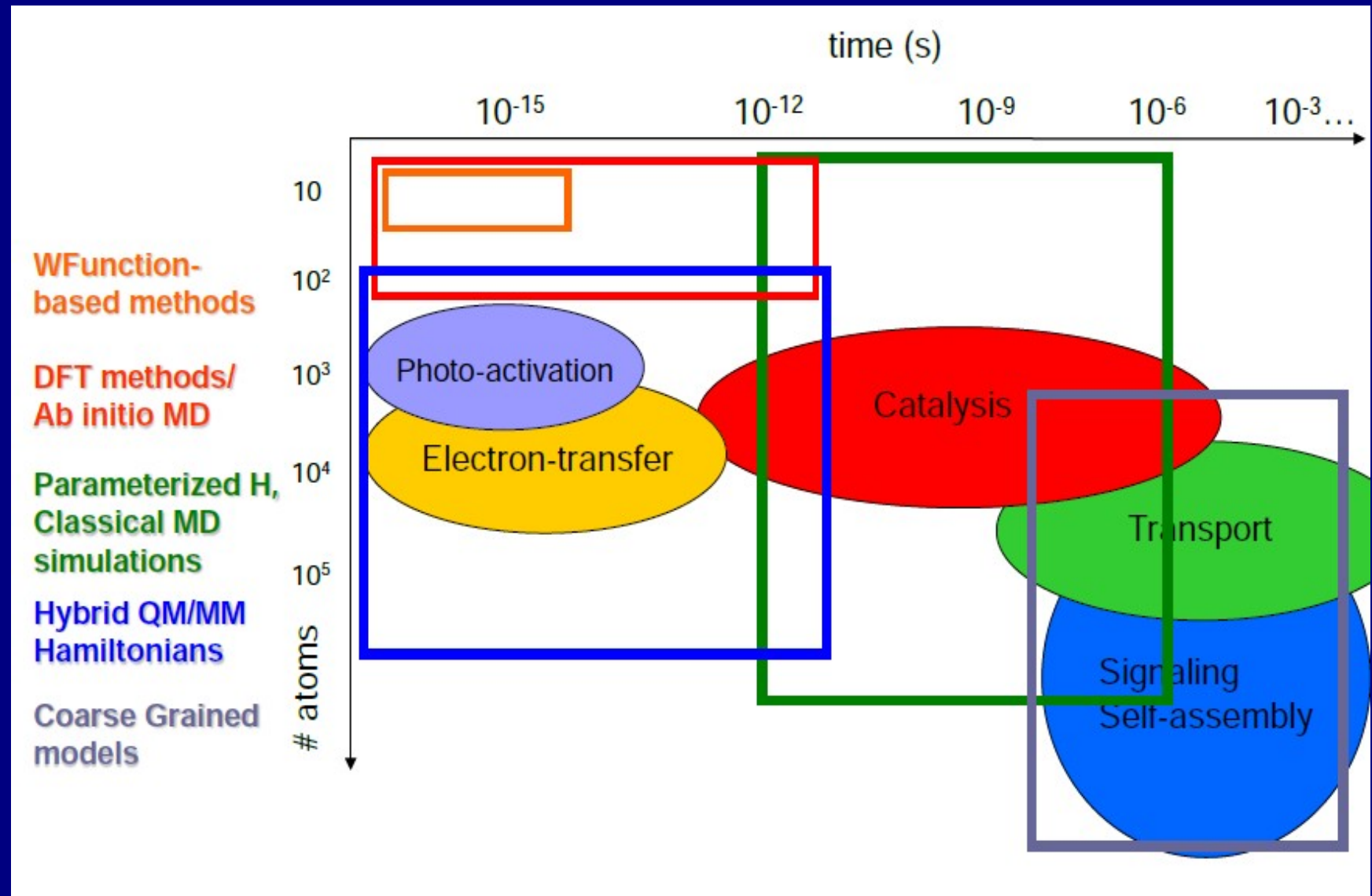
Conclusions



Molecular Scale & Complexity



Time and size scales



Theoretical Framework

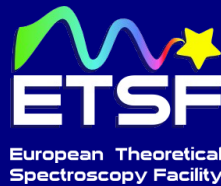


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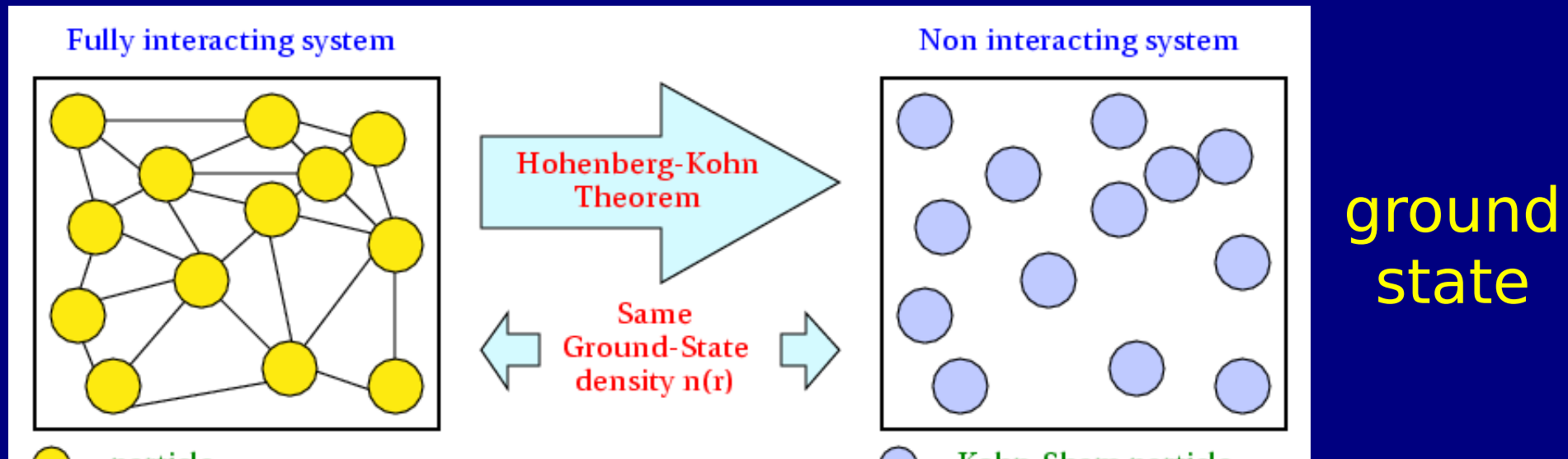
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Density Functional versus Many-body perturbation theory



What about excited-state properties: electron-ion dynamics, spectroscopies

DFT Success “~chemical accuracy”



EXACT ALTERNATIVE:

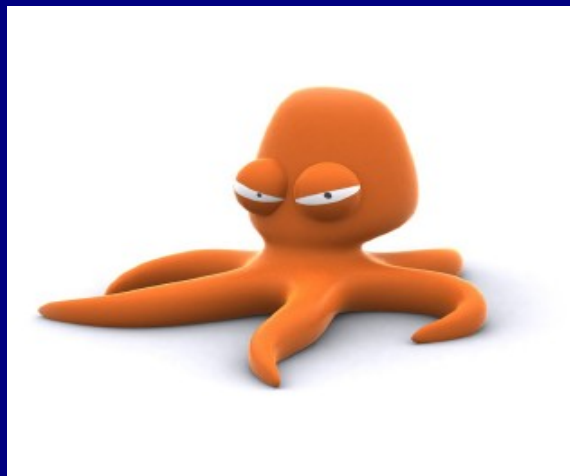
Time Dependent Density Functional Theory

All observables are functionals of the TD density

One-to-one correspondence between the time-dependent density and the external potential, $v(r,t) \longleftrightarrow \rho(r,t)$

$$i\hbar \frac{d}{dt} \Phi = H \Phi \quad \rightarrow \quad i\hbar \frac{d}{dt} \psi_i = H_{KS}[\{\psi_j\}] \psi_i, \quad i=1, \dots, N$$

$$H_{KS} = \frac{\hbar^2}{2m} (i\nabla - \frac{e}{c\hbar} (A + A_{xc}))^2 + V_{external} + V_{hartree} + V_{exchange} + V_{correlation}$$

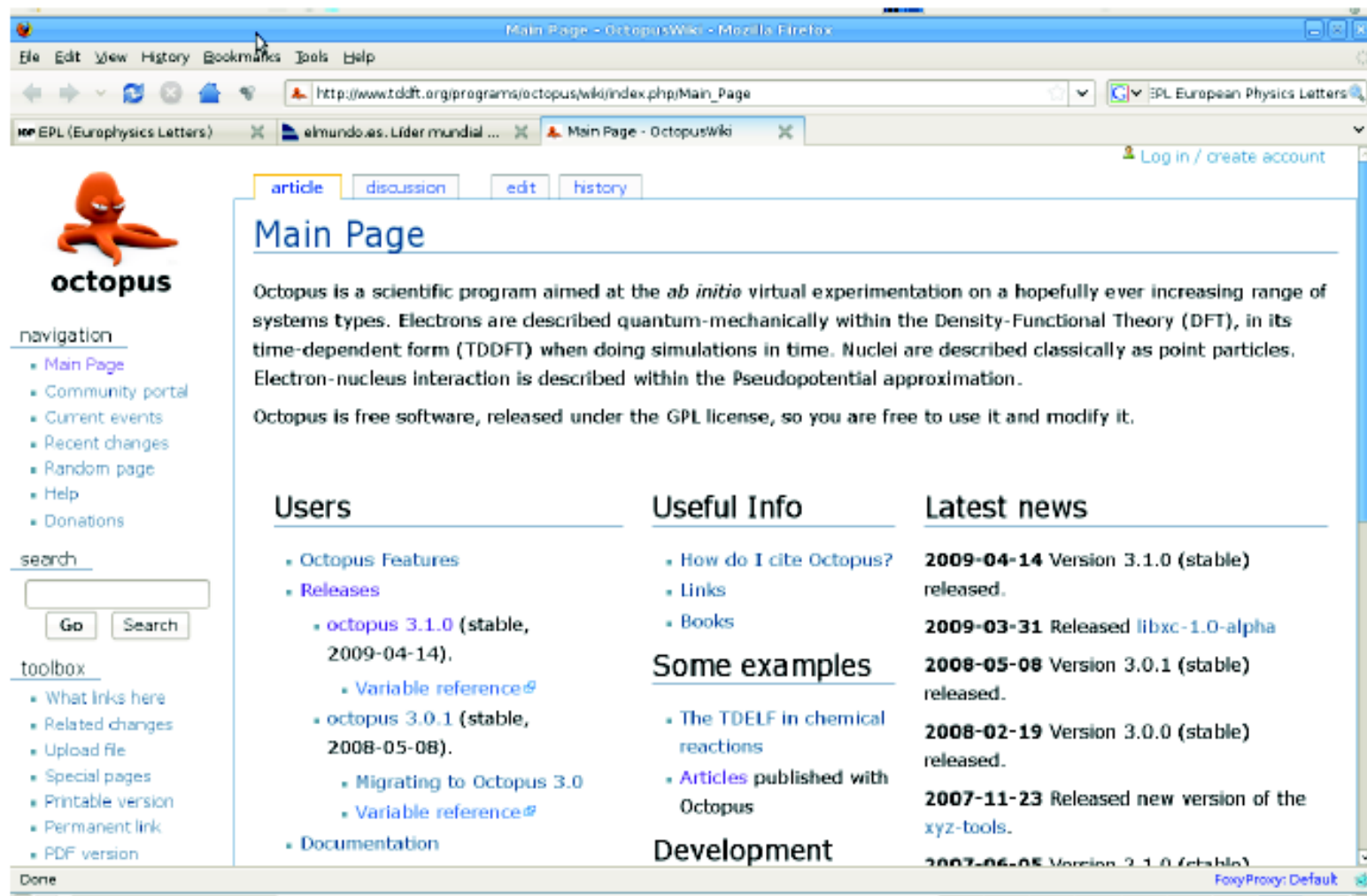


Linear and non linear phenomena accessible

Octopus Code <http://www.tddft.org>

Time-Dependent Density Functional Theory, Lecture Notes in Physics,
Vol. 706 (Springer, Berlin, 2006)





<http://www.tddft.org/programs/octopus/>



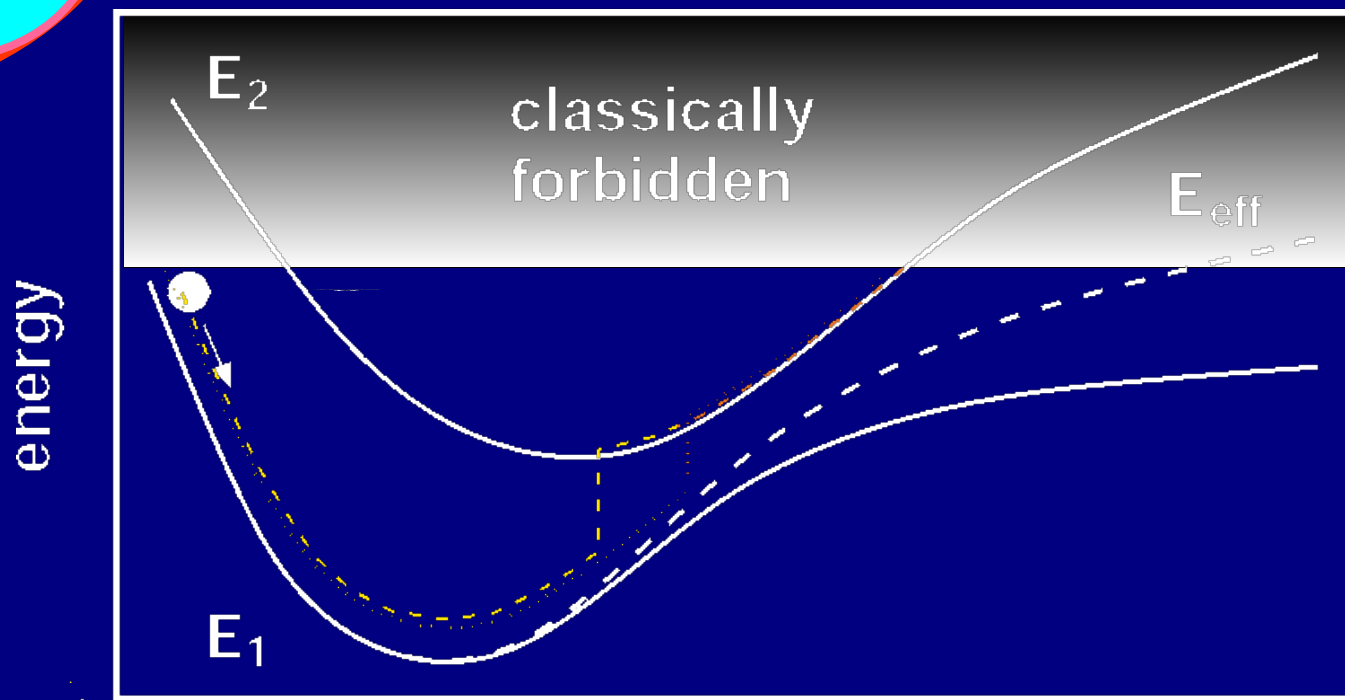
Ehrenfest's dynamics (non-adiabatic)

Electron dynamics

$$i\frac{\partial\Psi}{\partial t}=H\Psi \quad H\psi=E\psi$$

MD $F = m\alpha$

Multicomponent (TD)DFT: work in progress, way to describe the non-adiabatic processes



Surface hopping (Tully and Preston)

See review by N.L Doltsinis and D. Marx, J.Theo.Comp.Chem 1, 319 (2002)
and the book on TDDFT Springer Lecture Notes in Physics Vol 706 (2006)



TDDFT formulation for BO-MD dynamics: fictitious dynamics as in Car-Parrinello dynamics

The equation of motion without external field for simplicity

$$i\mu \frac{\partial}{\partial t} \psi_i = \left[\frac{-\hbar^2}{2m} \nabla^2 + V_{ion} + V_H + V_{xc} \right] \psi_i \quad M_I \ddot{R}_I = -\nabla_I E[\psi, R]$$

J.L. Alonso, X. Andrade, P. Echenique, F. Falceto, D. Prada, AR PRL (2008), JCTC (2009), NJP (2010)



Physical interpretation

Electronic time

$$\mu \frac{d}{dt} = \frac{d}{d(t/\mu)}$$

- Two time scales: real for ions and fictitious for electrons.

Maximum time step

$$\Delta t_{max}(\mu) = \mu \Delta t_{max}(\mu = 1)$$

Scaling of electronic excitation energies

$$\omega_i(\mu) = \frac{1}{\mu} \omega_i(\mu = 1)$$



TDDFT Lagrangian formulation for BO-MD dynamics: fictitious dynamics: performance

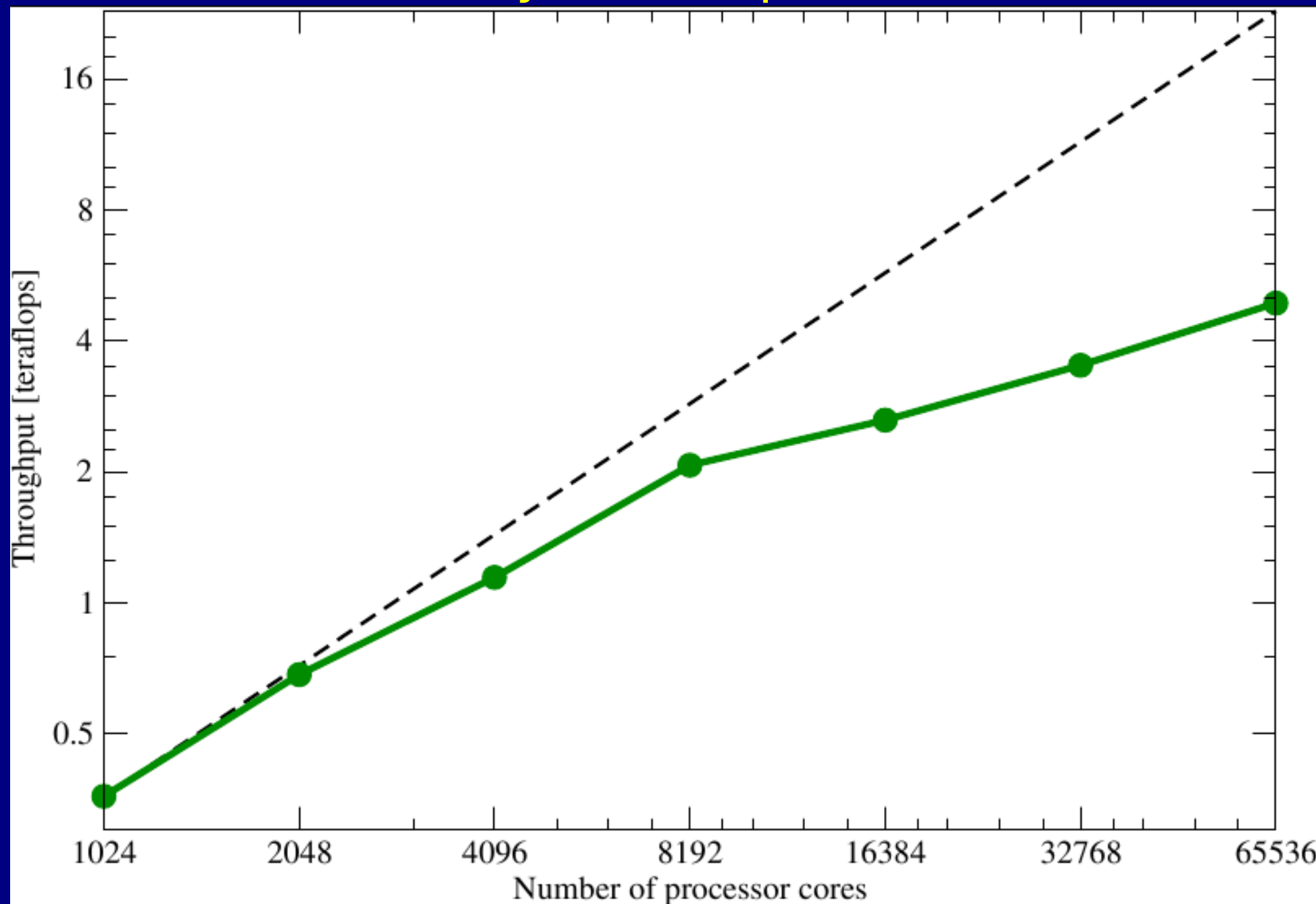


Illustration: Some applications

Optical properties of nanostructures and biomolecules
(excite state dynamics and energy store/harvesting)
Molecular level alignment

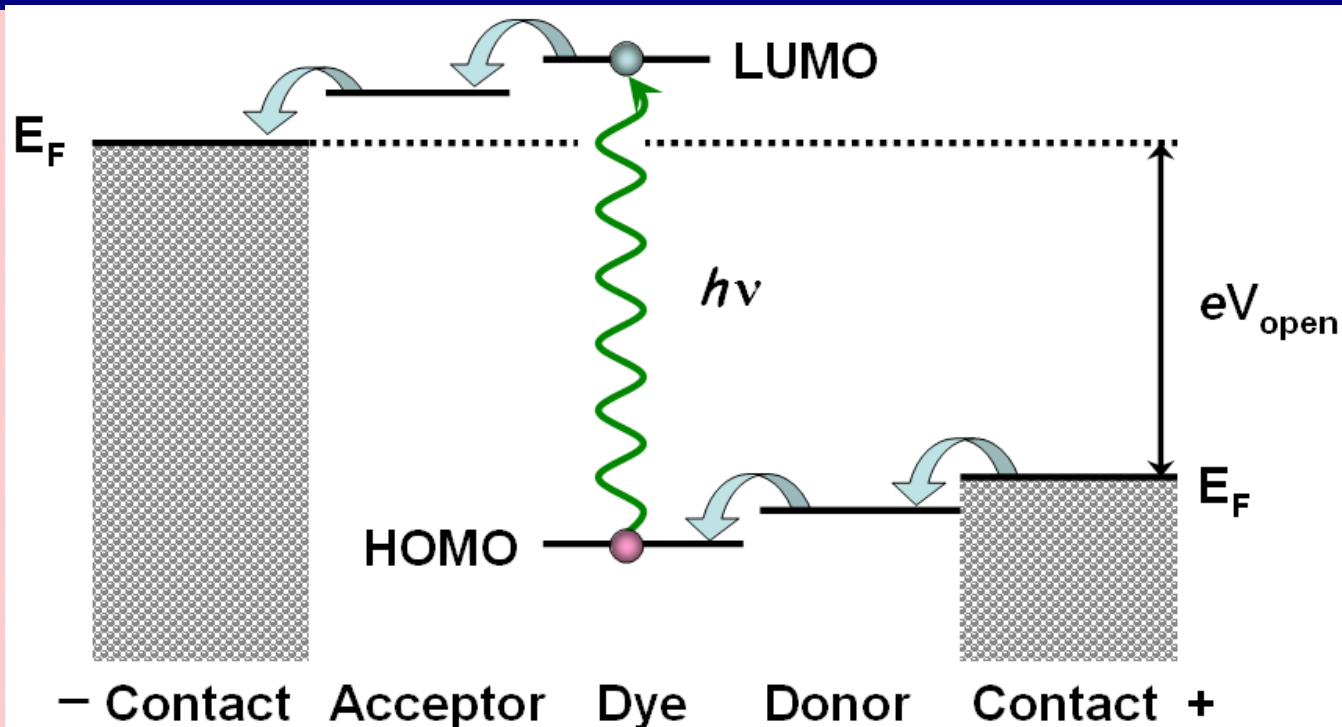
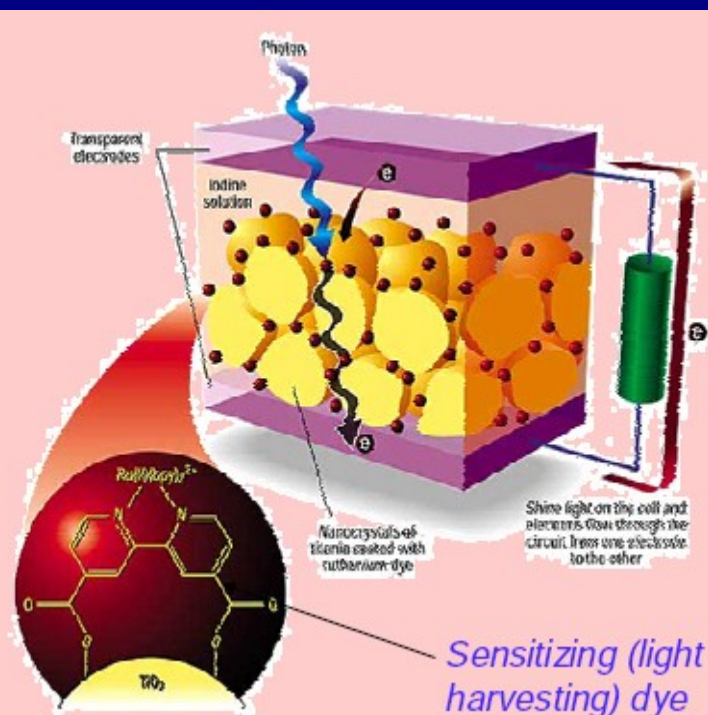
Optimal control theory: High-harmonic generation

Nanosolvation????



Charge Transfer excitations

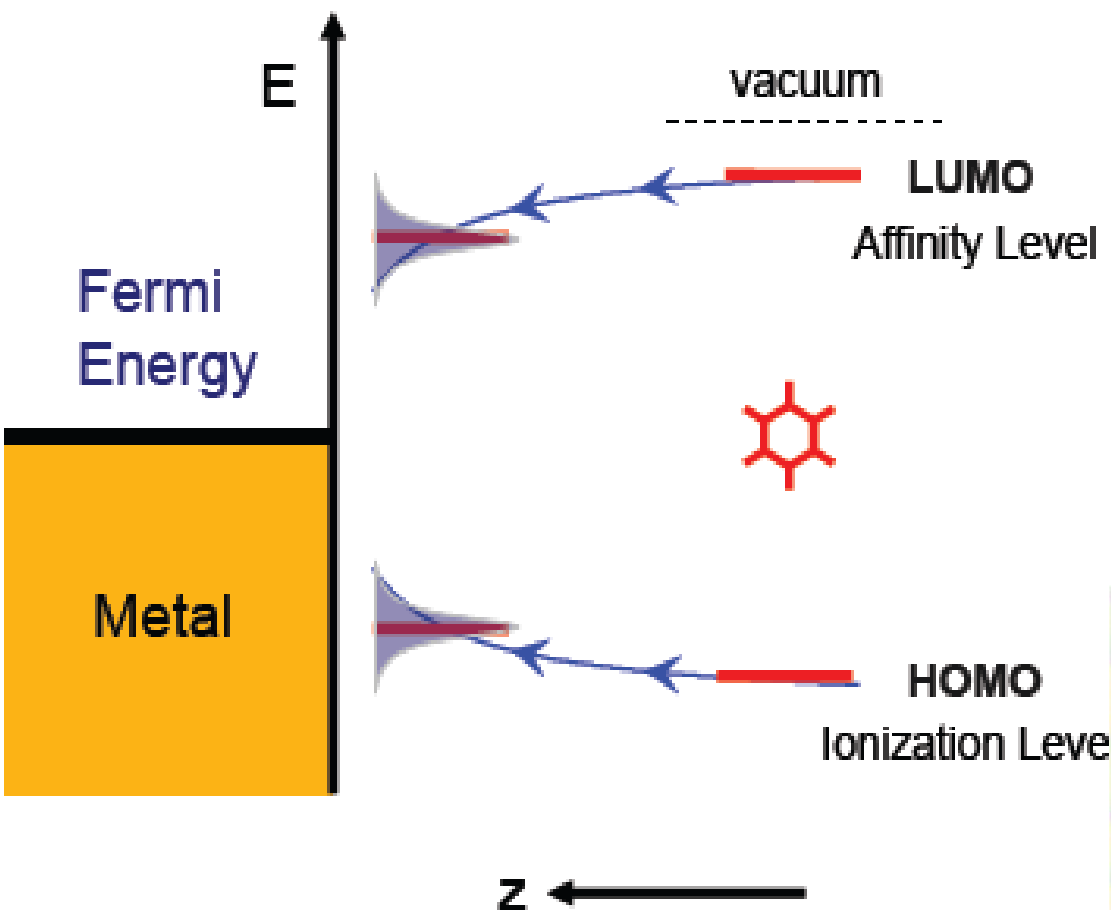
Photovoltaic Hybrids: Grätzell cells



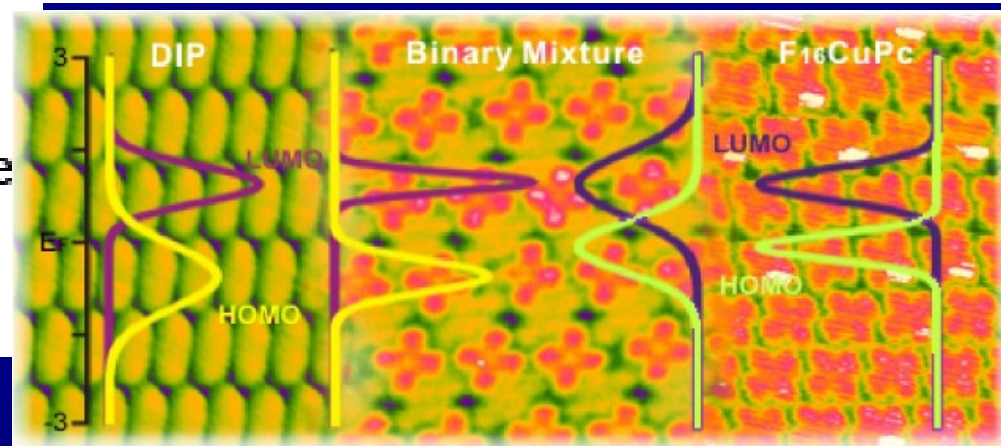
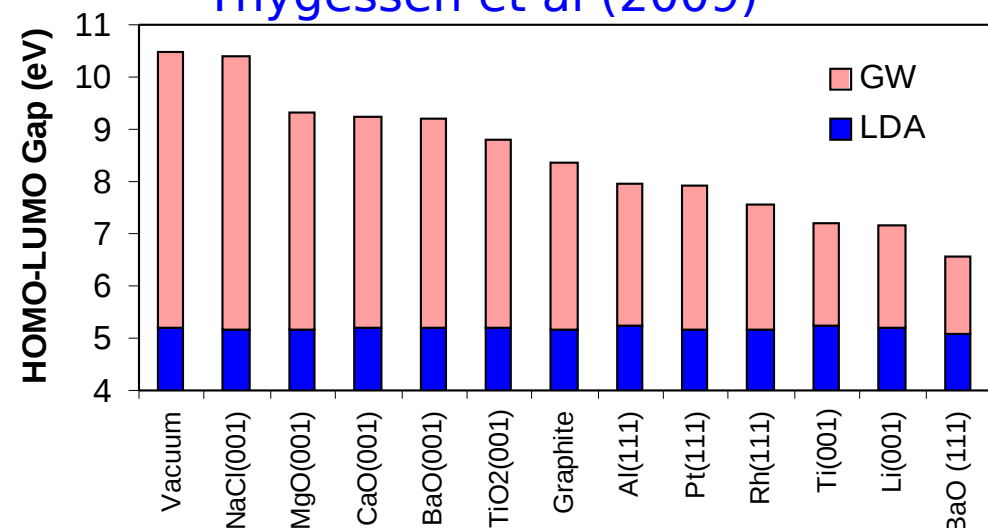
Road to Improving Efficiency: Controlling position of the frontier orbitals



LEVEL ALIGNMENT, Dispersion forces VdW



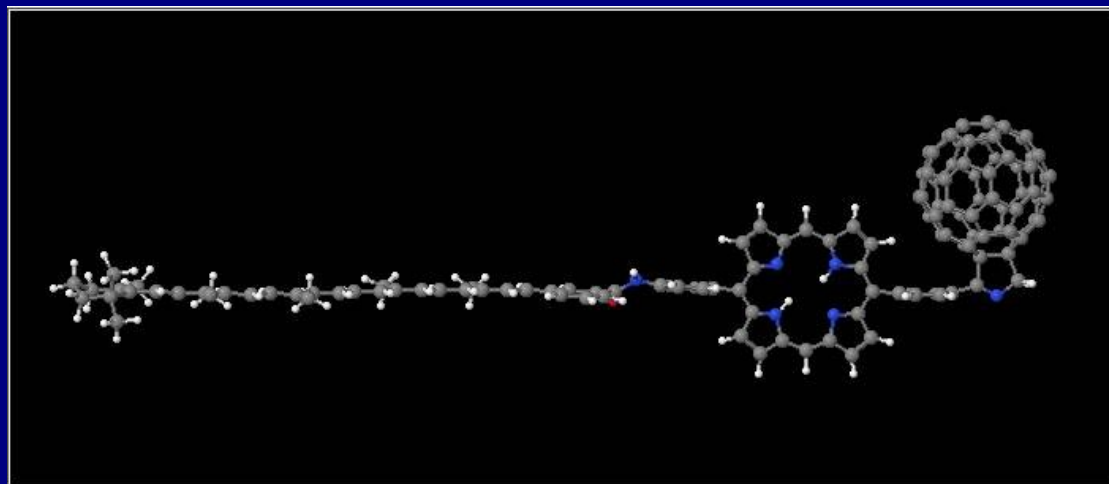
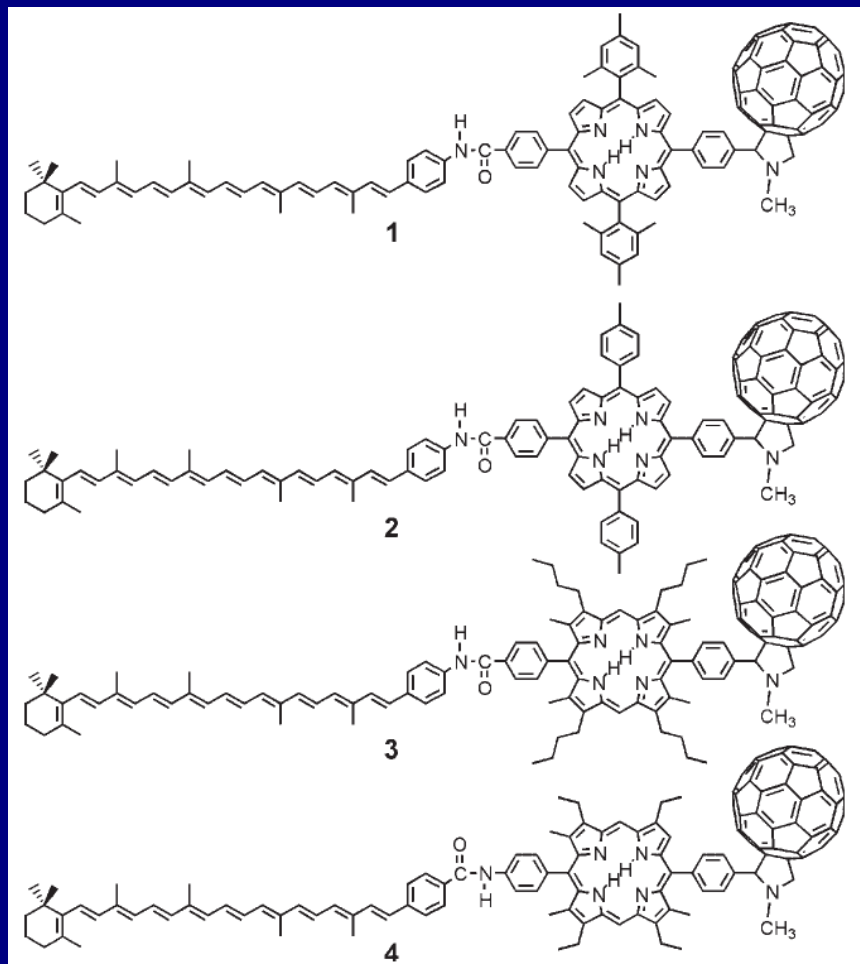
Thygesen et al (2009)



Supramolecular chemistry, Biophysics , Molecular Transport.....



Photo-excitation of the light-harvesting carotenoid-porphyrin-C60



Comparing HOMO and LUMO energies

molecules	ϵ_{HOMO} (eV)	ϵ_{LUMO} (eV)
β -carotenoid	-4.47	-3.49
Diaryl-porphyrin	-5.15	-3.31
pyrrole-C ₆₀	-5.89	-4.42
Triad	-4.51	-4.38



Photo-excitation of the light-harvesting carotenoid-porphyrin-C60

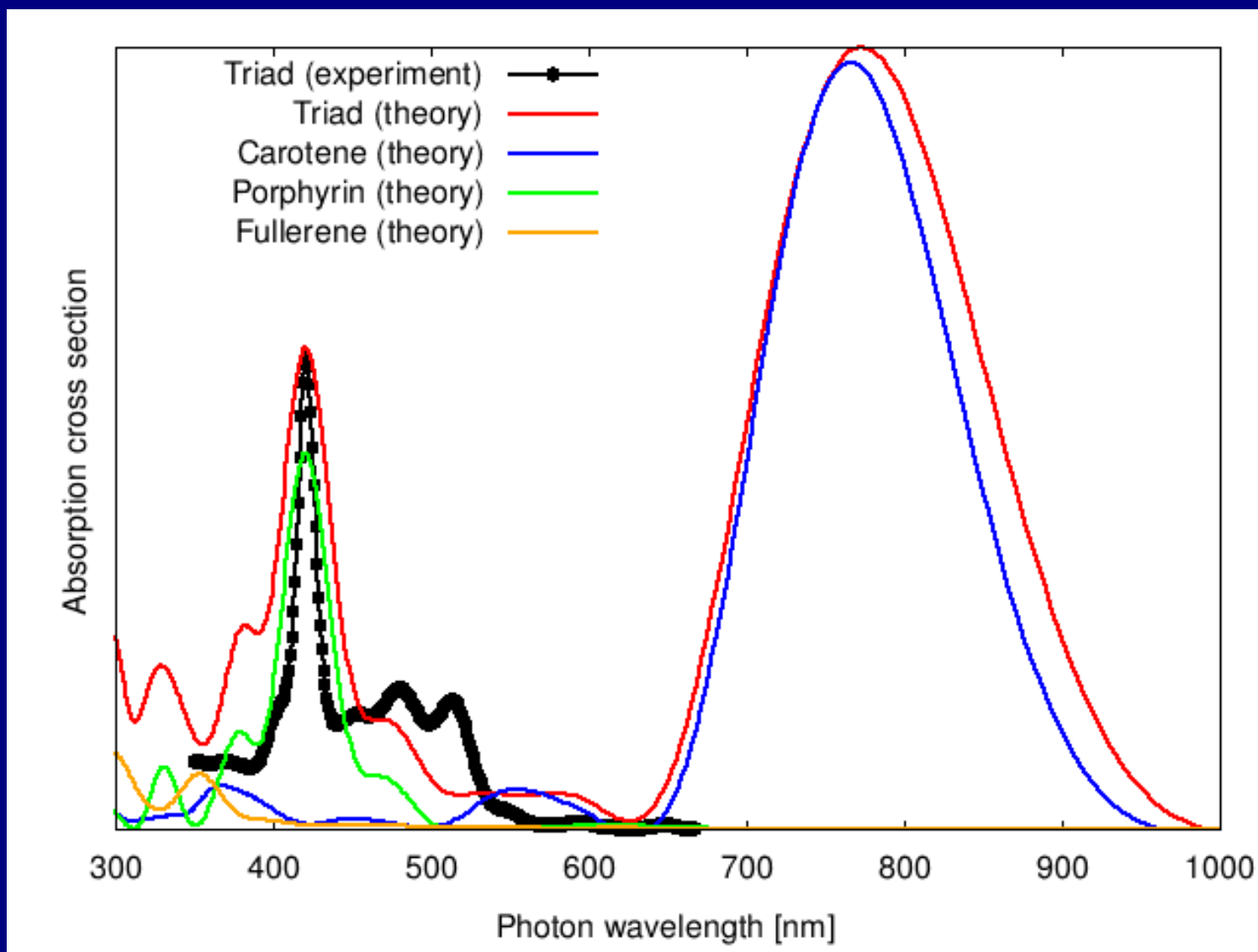
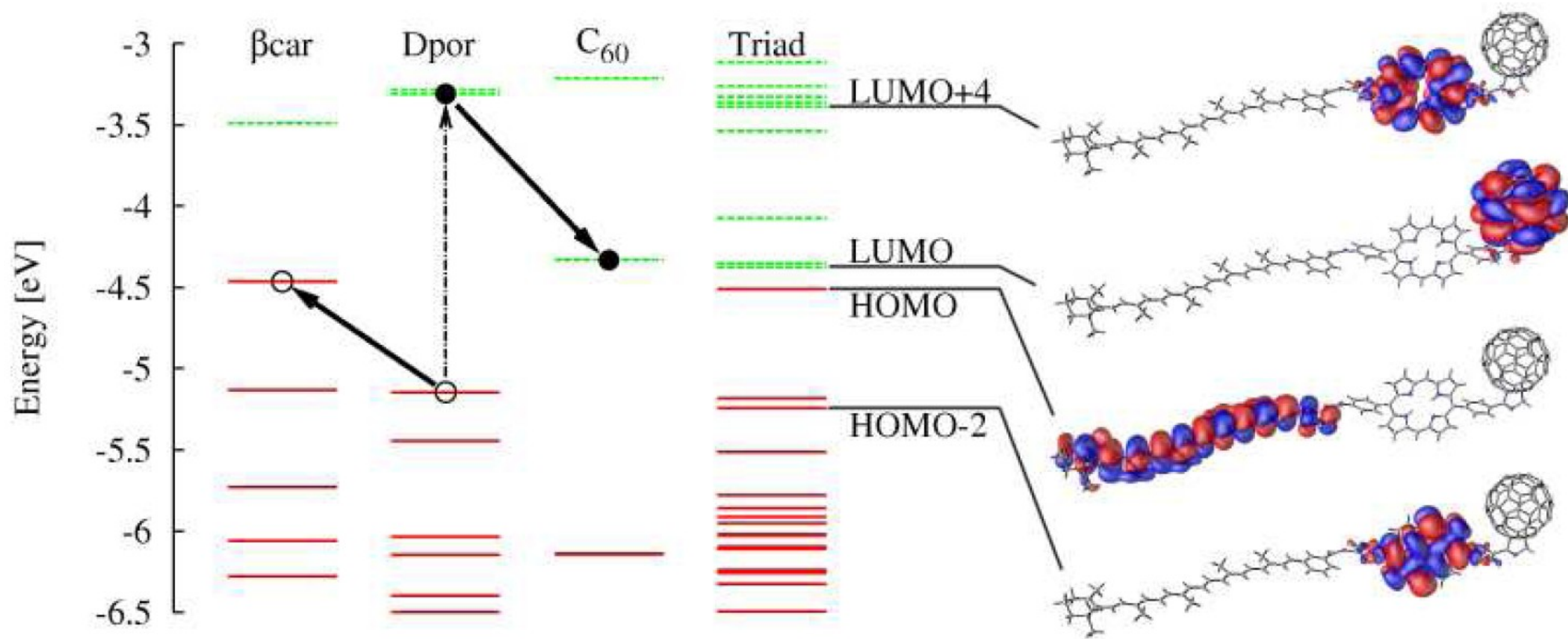
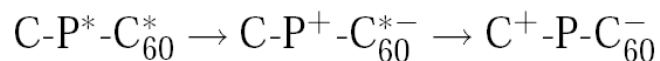
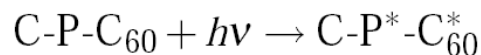
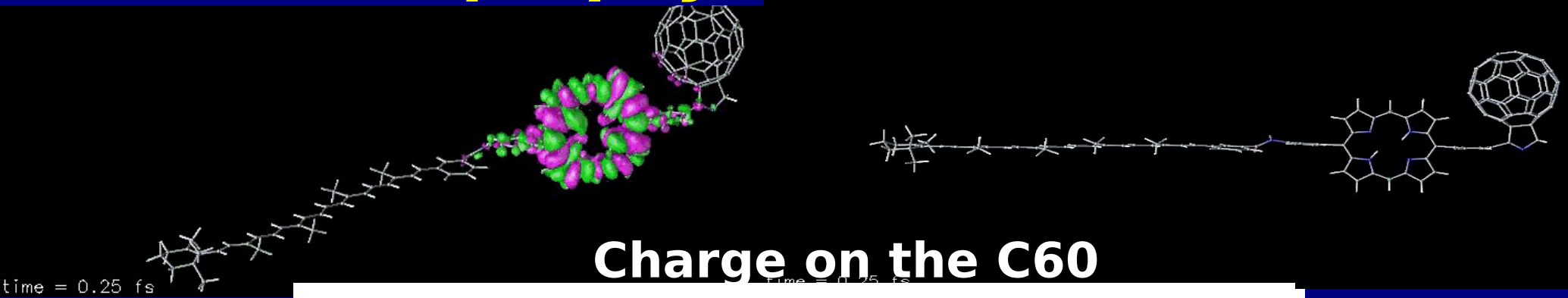


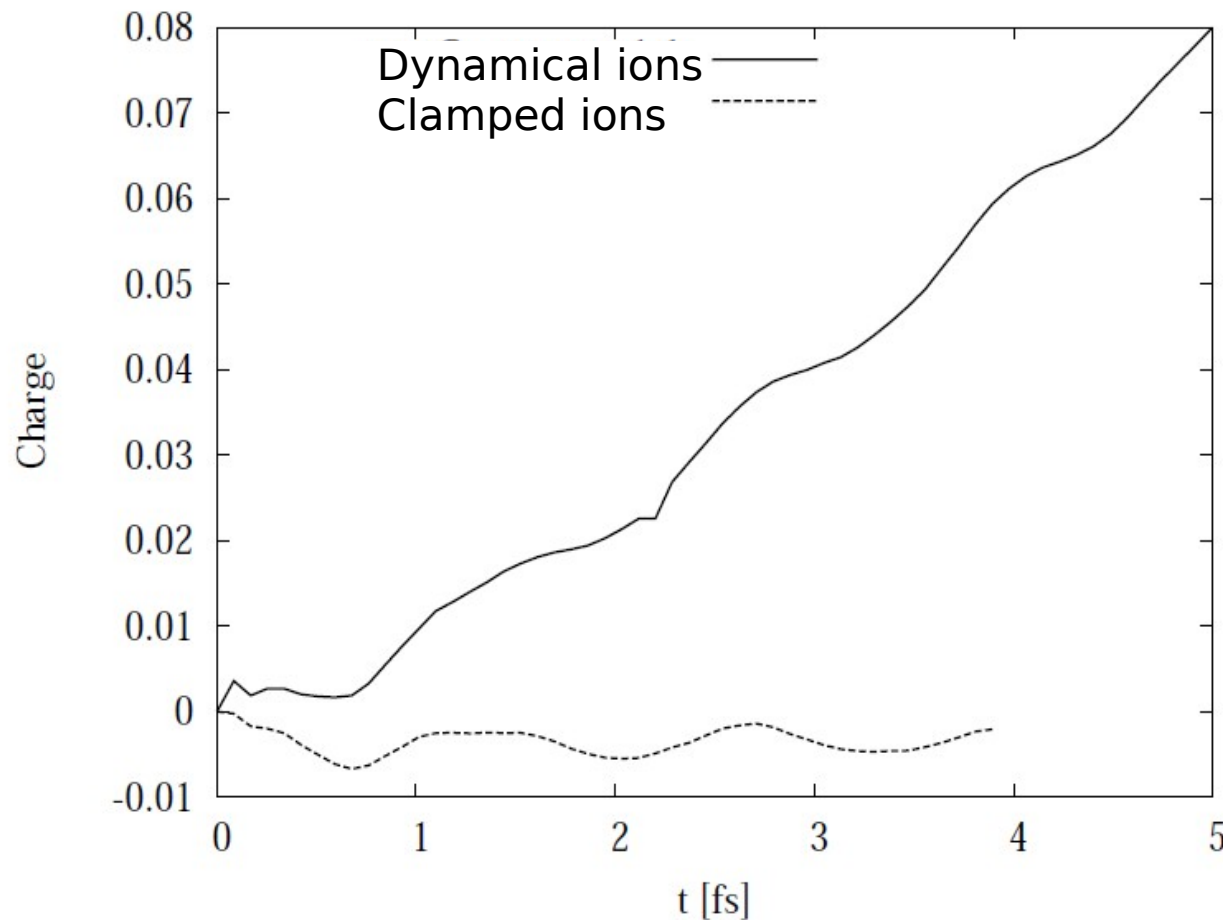
Photo-excitation of the light-harvesting carotenoid-porphyrin-C60

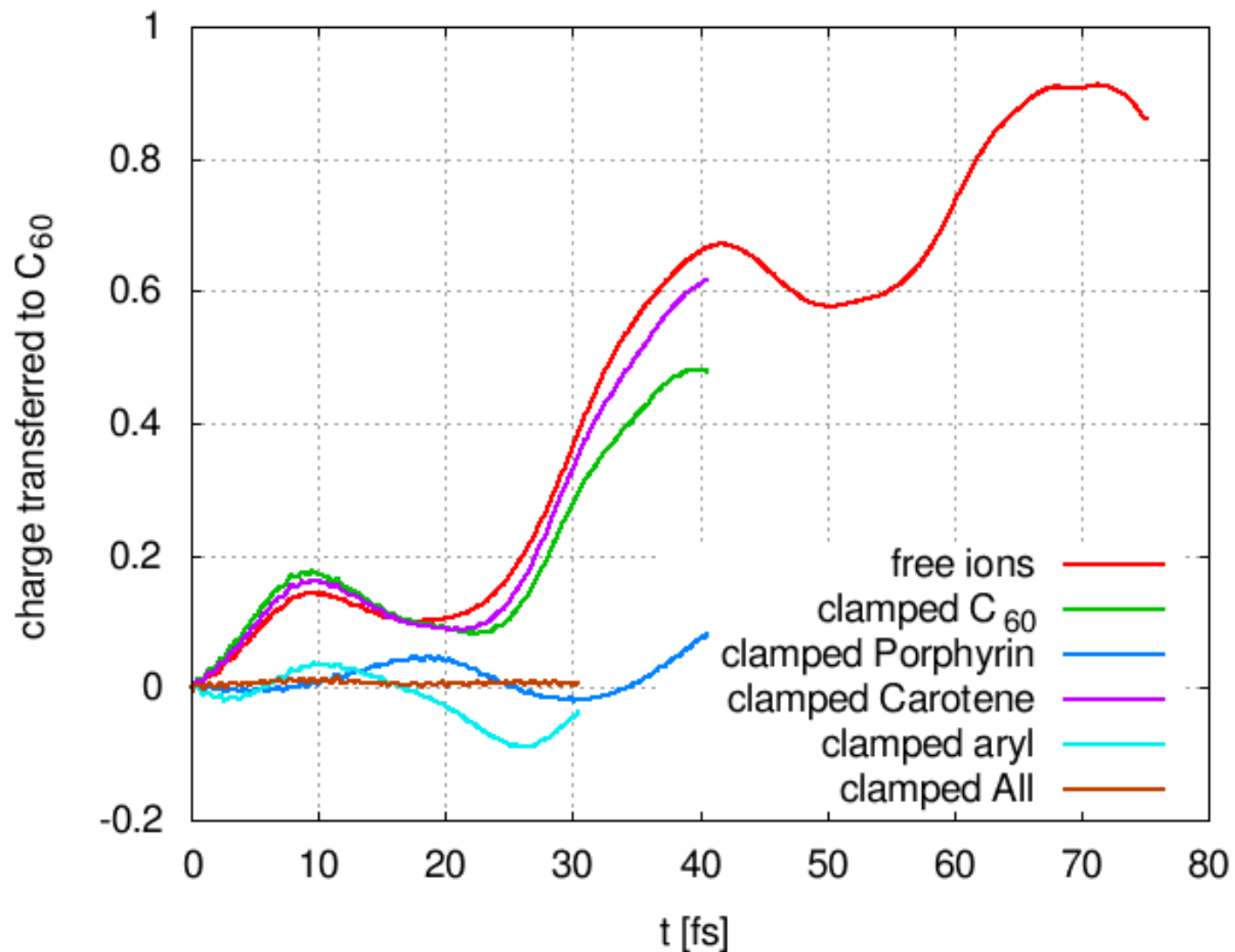


Carotenoid-porphyrin-C60 photodynamics

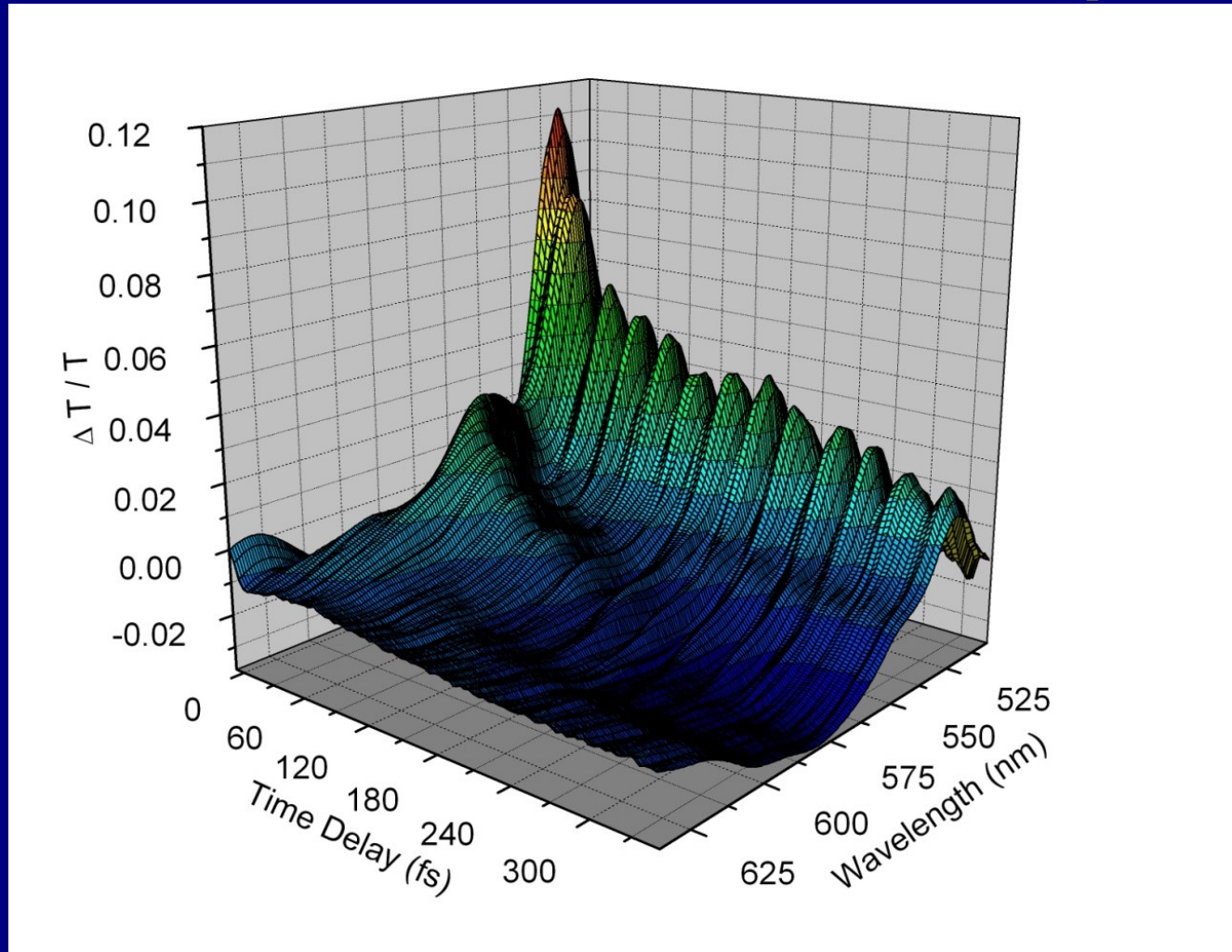


Charge on the C60





Simulation of the Time resolved spectra

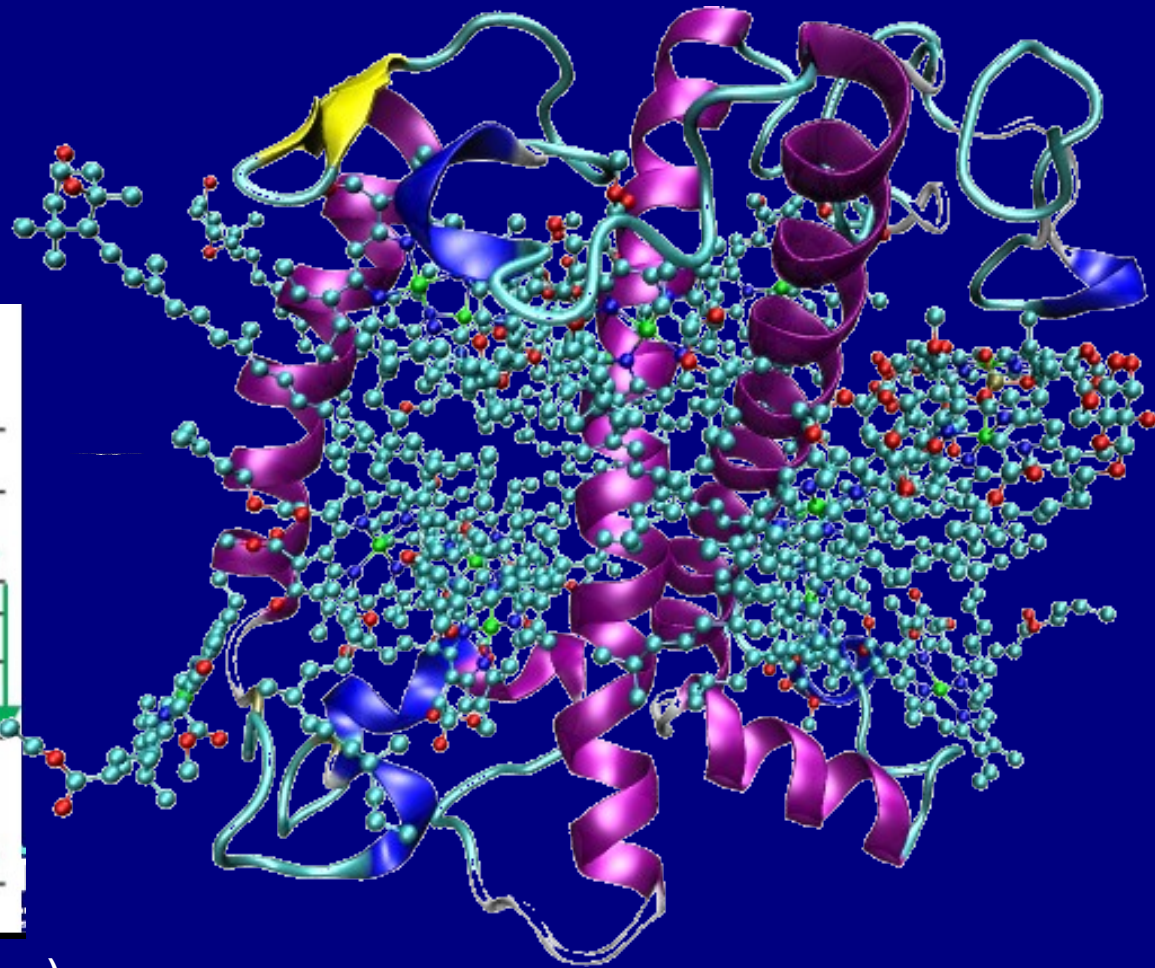
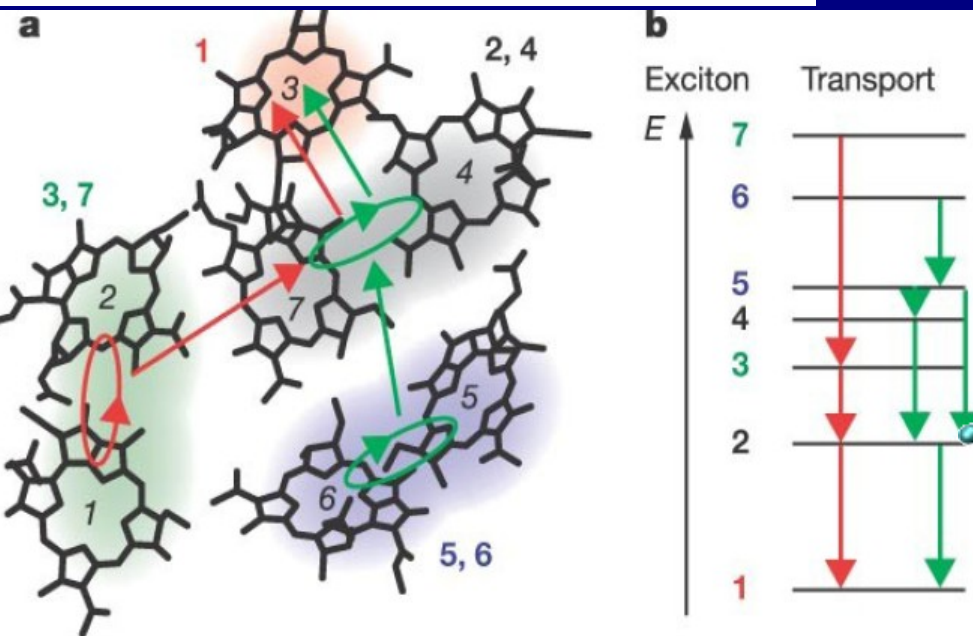
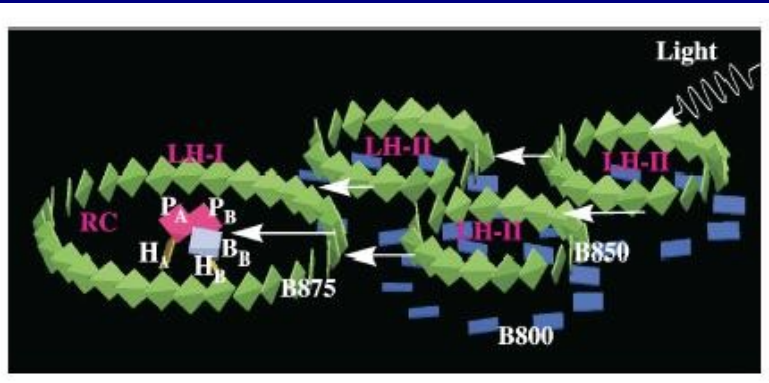


The carotene feature around 523 nm oscillates both in intensity and energy, the porphyrin feature red-shifts from about 540 nm to 573 nm and decays almost completely within the first 100 fs after photoexcitation



Towards the color of plants:

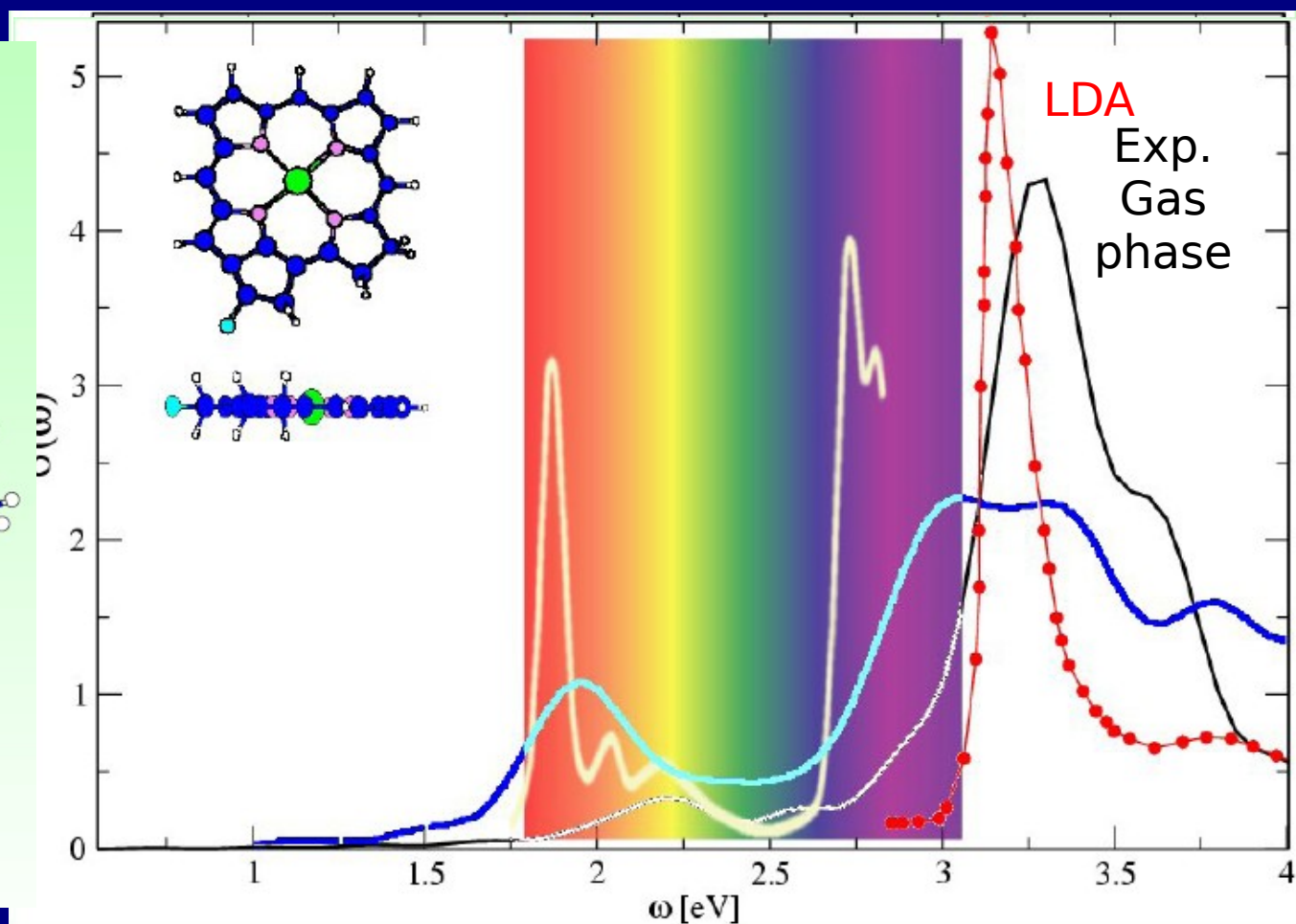
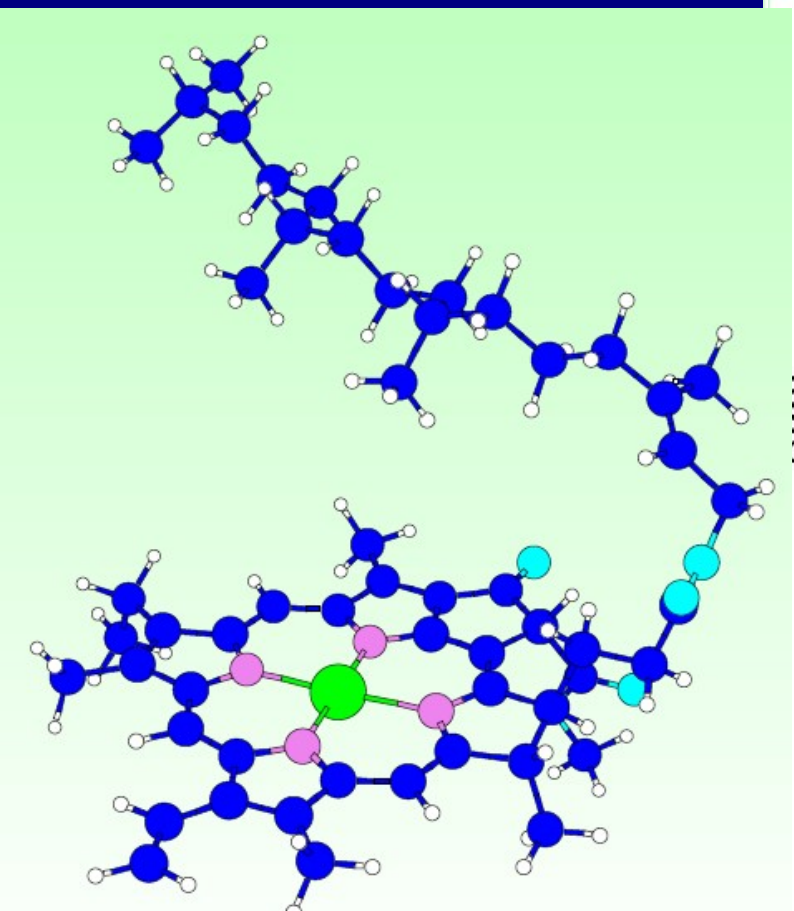
e.g. Spinach



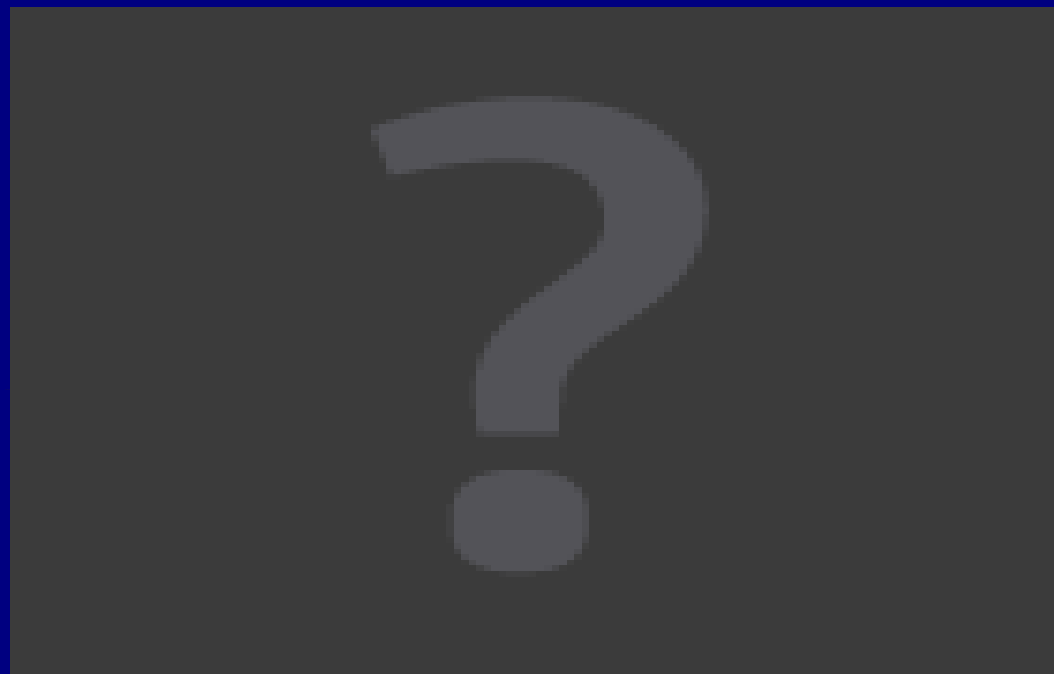
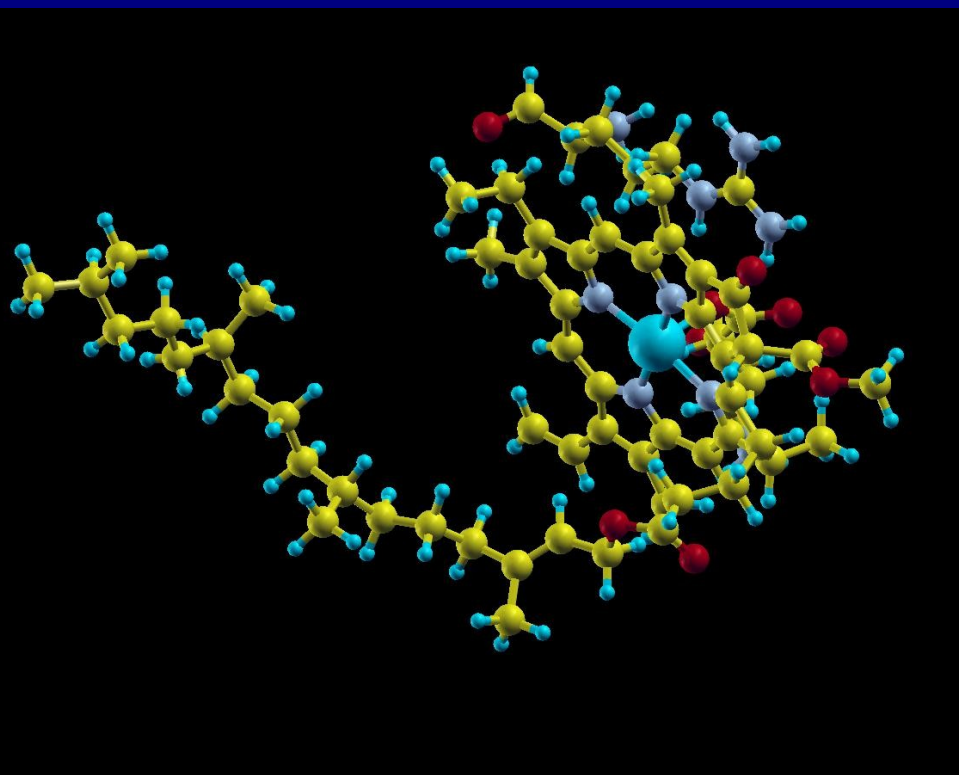
Fleming et al Nature (2005),.... (Berkeley)



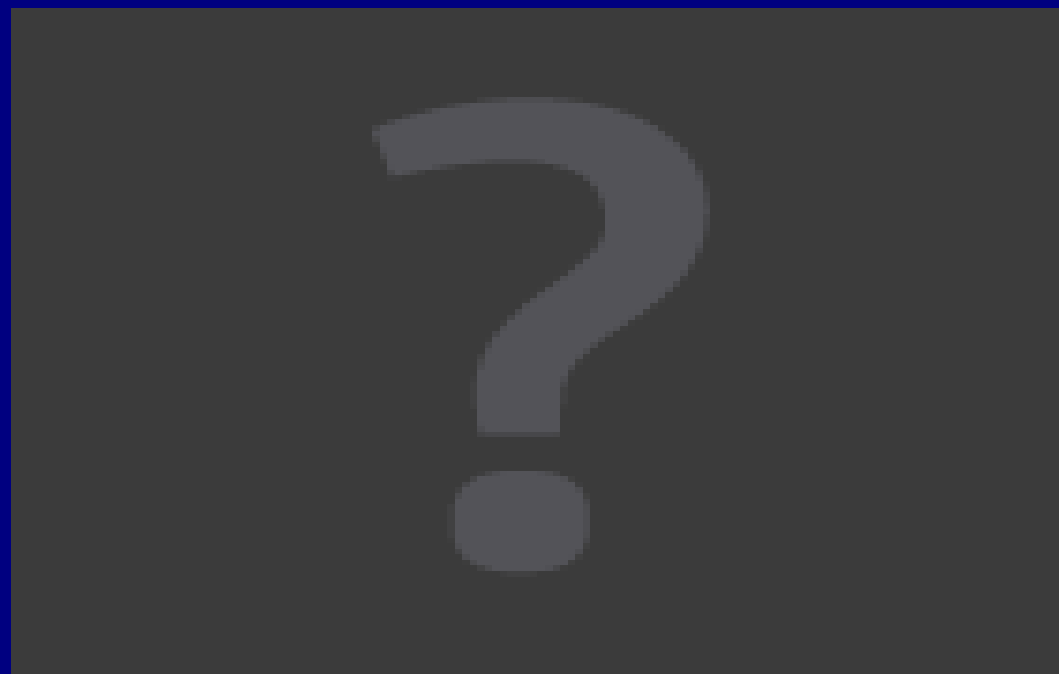
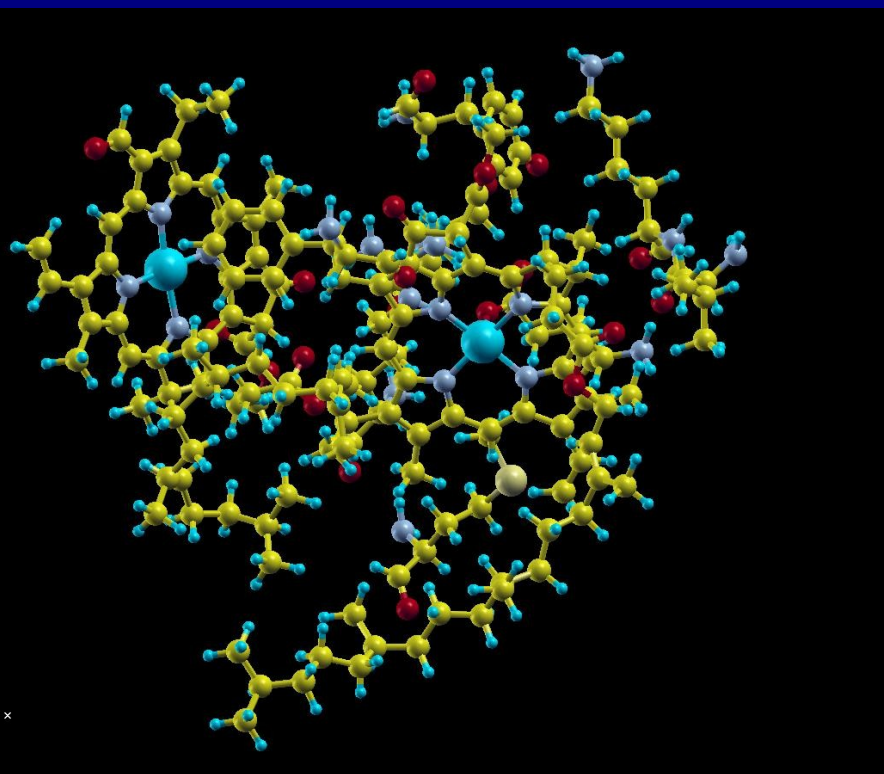
Towards the color of plants: The real Chlorophyll molecule



Towards the color of plants:



Towards the color of plants:



Challenge: proper treatment of the whole environment, including dissipation for very long simulations from fs to milliseconds and more.....



Optimal control theory

Key question: What is the laser pulse that drives the system into a predefined goal?

Procedure: Define a target operator \hat{O} and at the end of the laser interaction ($t = T$) maximize the functional

$$J_1[\Psi] = \langle \Psi(T) | \hat{O} | \Psi(T) \rangle = |\langle \Psi(T) | \Phi_F \rangle|^2$$

$\hat{O} = |\Phi_F\rangle\langle\Phi_F|$ target state

W. Zhu, J. Botina, H. Rabitz, JCP 108, 1953 (1998)

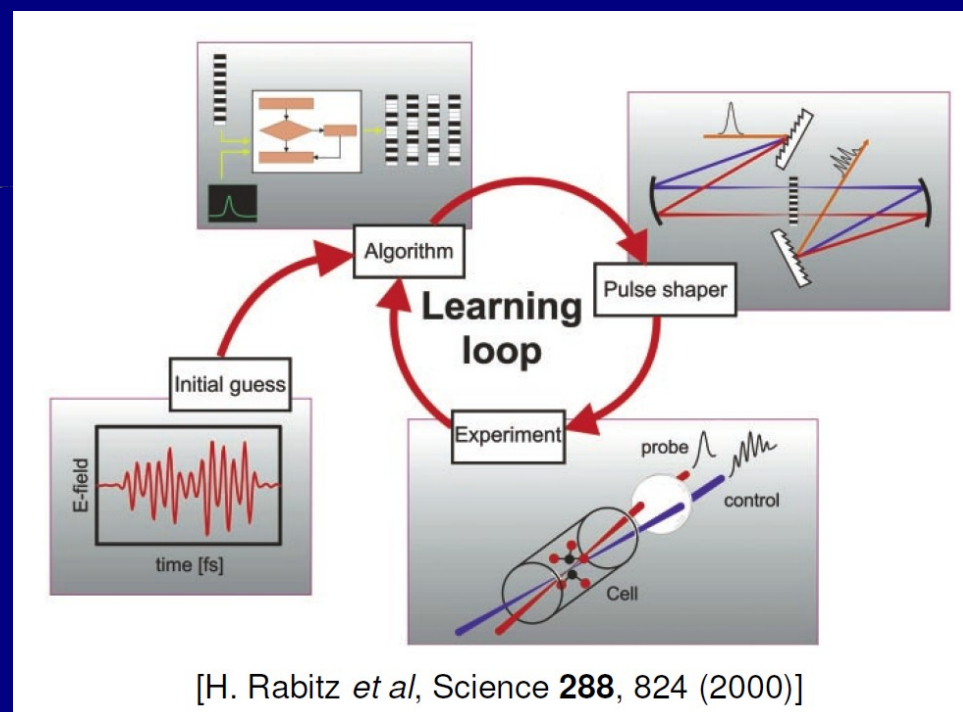


What makes experimental “control” possible

- Existence of laser sources, since the 1960’s.
- Femto-second laser sources, which allow for fast processes (avoiding decoherence), and extending the band-width.
- High-intensities.
- Laser shapers.
- Learning-loops algorithms.

Applications:

Selective bond-breaking
High-harmonic generations



Design of bond-breaking laser pulses

THE JOURNAL OF CHEMICAL PHYSICS **127**, 201101 (2007)

Coherent control of bond breaking in amino acid complexes with tailored femtosecond pulses

T. Laarmann,^{a)} I. Shchatsinin, P. Singh, and N. Zhavoronkov

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(Received 6 August 2007; accepted 15 October 2007; published online 30 November 2007)

Intense femtosecond laser pulses, judiciously tailored in an adaptive, optimal control feedback loop were used to break preferentially the acyl-N ("peptide") bond of Ac-Phe-NHMe that may be regarded as a dipeptide model. We show that coherent excitation of complex wave packets in the strong-field regime allows to cleave strong backbone bonds in the molecular system preferentially, while keeping other more labile bonds intact. These results show the potential of pulse shaping as a powerful complementary analytical tool for protein sequencing of large biopolymers in addition to the well-known mass spectrometry and chemical analysis. © 2007 American Institute of Physics.

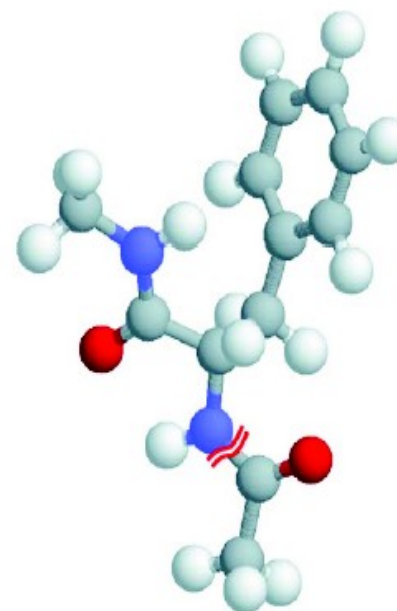


FIG. 1. (Color online) Optimal tailoring of intense femtosecond light can be used to preferentially break peptide bonds, such as the indicated N1-C3 bond in the amino acid complex Ac-Phe-NHMe.

K. Kriger, A. Castro, E.K.U. Gross, (unpublished)



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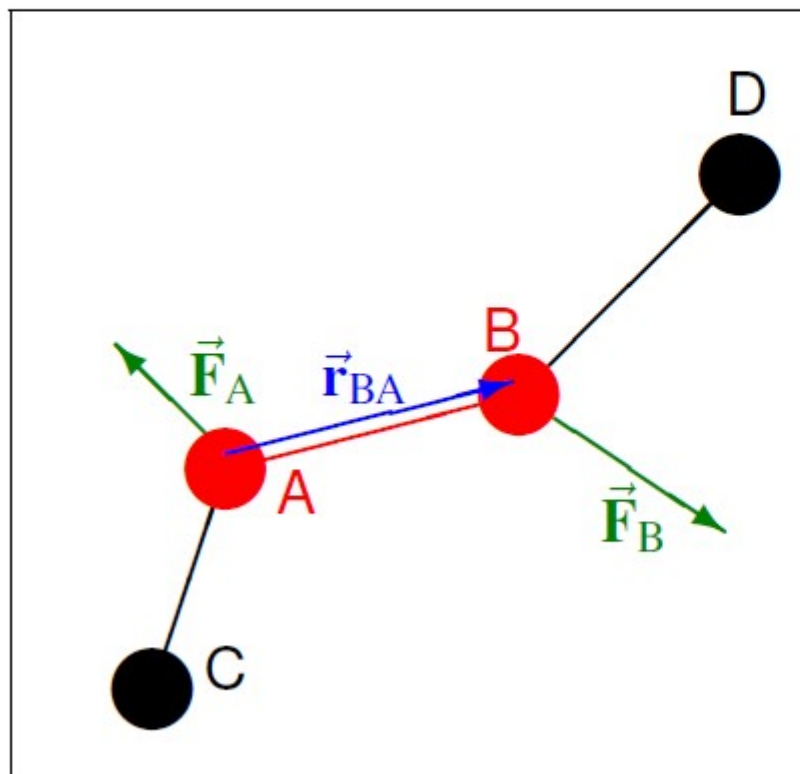
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Design of bond-breaking laser pulses (II)



- If the AB bond is breaking,

$$\vec{r}_{BA} \cdot \vec{F}_B > 0$$

$$\vec{r}_{BA} \cdot \vec{F}_A < 0$$

\Rightarrow

$$\vec{r}_{BA} \cdot \vec{F}_B - \vec{r}_{BA} \cdot \vec{F}_A > 0 \Rightarrow$$

$$\vec{r}_{BA} \cdot \vec{F}_{BA} > 0$$

$$(\vec{F}_{BA} = \vec{F}_B - \vec{F}_A)$$

Design of bond-breaking laser pulses (III)

- Bond-breaking target:

$$J_1[\Psi] = \frac{1}{t_f - t_0} \int_{t_0}^{t_f} dt \langle \Psi(t) | |\vec{\mathbf{r}}_{\text{BA}}| |\hat{\vec{\mathbf{F}}}_{\text{BA}}(t)| \left(1 + \vec{\mathbf{r}}_{\text{BA}} \cdot \hat{\vec{\mathbf{F}}}_{\text{BA}}(t) \right) | \Psi(t) \rangle .$$

- Within TDDFT, $\hat{\vec{\mathbf{F}}}_{\text{A}}$ is a local density operator: the necessary expectation values can be written as simple density integrals.

$$\hat{\vec{\mathbf{F}}}_{\text{A}}(t) = \int d^3r \hat{n}(\mathbf{r}) \nabla_{\vec{\mathbf{r}}_{\text{A}}} v_{\text{KS}}(\mathbf{r}, t) .$$

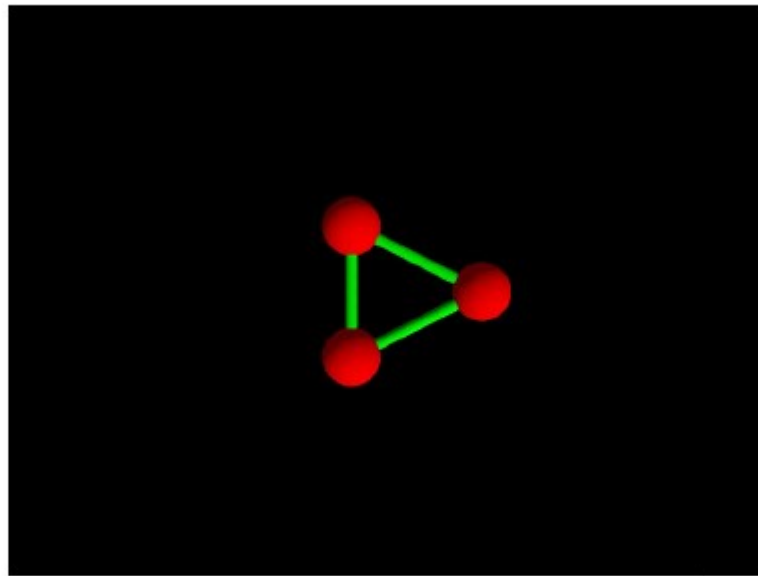
- During the time of the pulse, the nuclei are *fixed*. The photo-reaction occurs afterwards, driven by the excited electronic system.

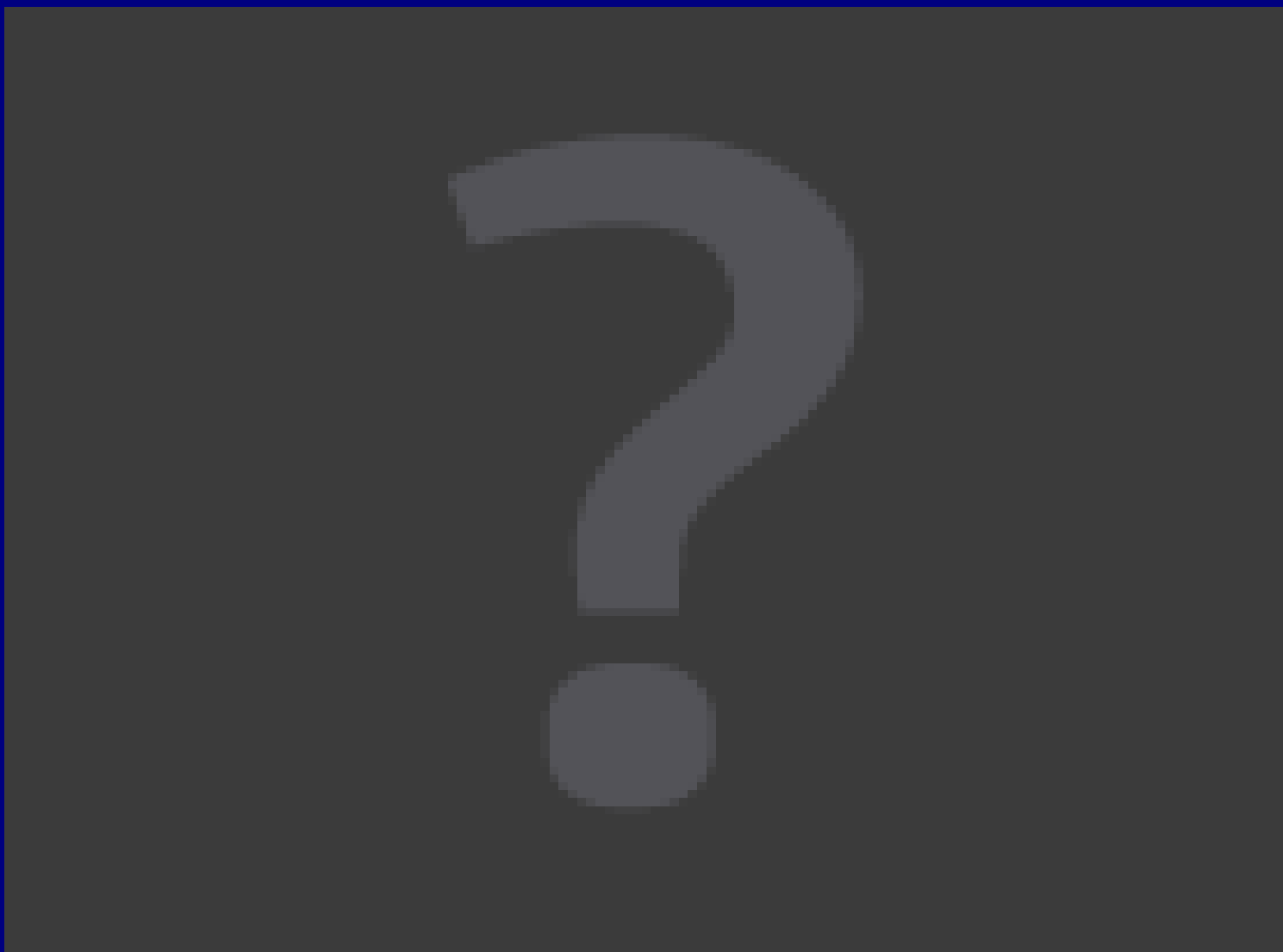


Design of bond-breaking laser pulses (IV)

- Case study: selective photo-dissociation of H_3^+ into H_2 and H^+ .
- For this case, we used a gradient-free optimization scheme, which only necessitates forward propagations.
- The field is expanded as a Fourier series:

$$\epsilon(t) = \sum_{n=1}^{N/2} a_n \sqrt{\frac{2}{T}} \cos\left(\frac{2\pi}{T} nt\right) + b_n \sqrt{\frac{2}{T}} \sin\left(\frac{2\pi}{T} nt\right)$$





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TDDFT is a powerful tool to “predictive” describe the combined dynamics of electron/ion in response to external electromagnetic fields of large scale nanostructures, biological molecules and extended systems spanning very different time scales

ETSF



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Halle

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Modena

K. Thygesen

Denmark



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Thank you!!!!

For more details see:
<http://nano-bio.ehu.es>
<http://etsf.eu>

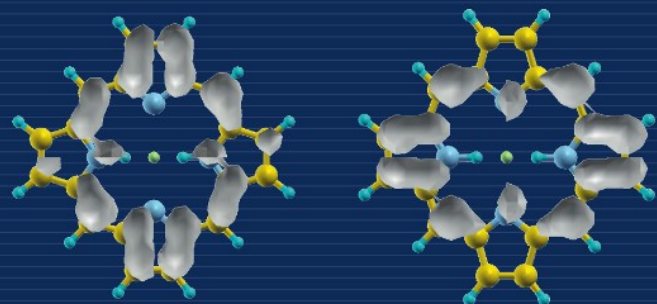


Free standing porphyrins

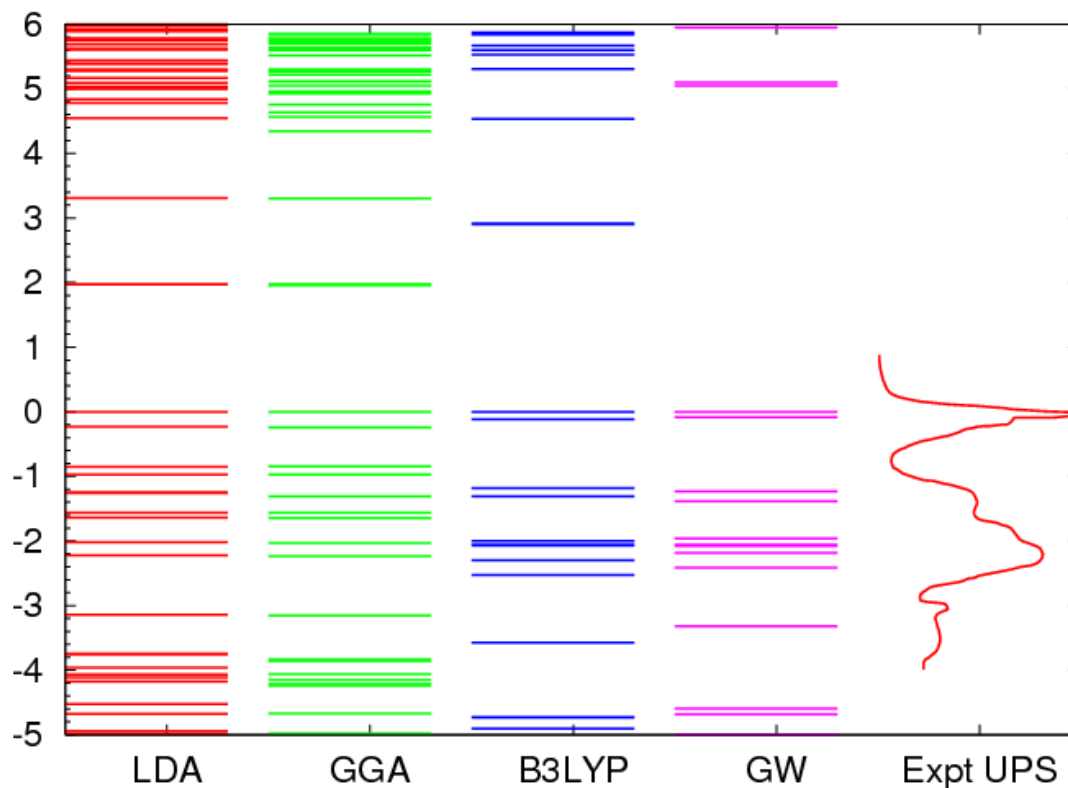
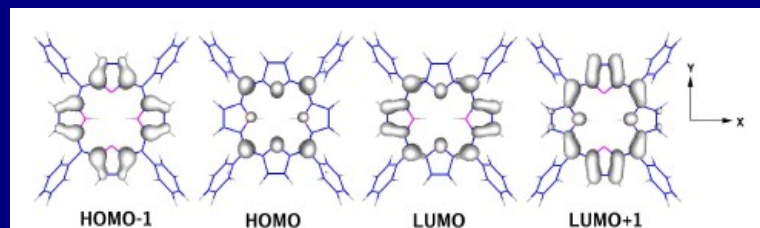
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Expt UPS: Ghosh, *Theor Chem Acc* **97**, 143 (1997); Dubois, *Chem Phys Lett* **75**, 434 (1980)

Journal of Chemical Physics 131, 84101 (2009)



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Nano solvation

Spectroscopic Probe for Characterizing Solvent-Mediated Folding in Dicarboxylate Dianions**
Torsten Wende, Marius Wanko, Ling Jiang, Gerard Meijer, Knut R. Asmis and AR (2010)



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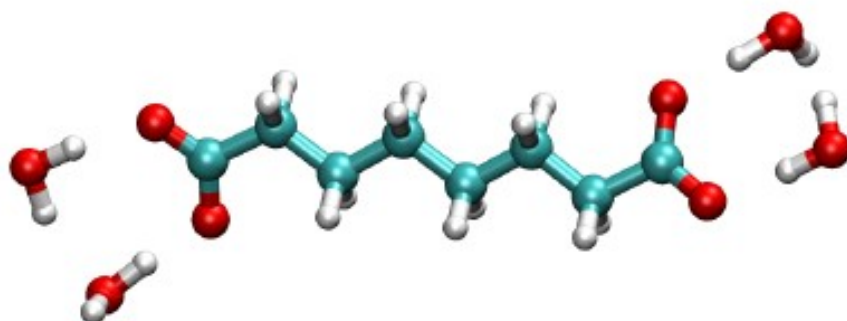


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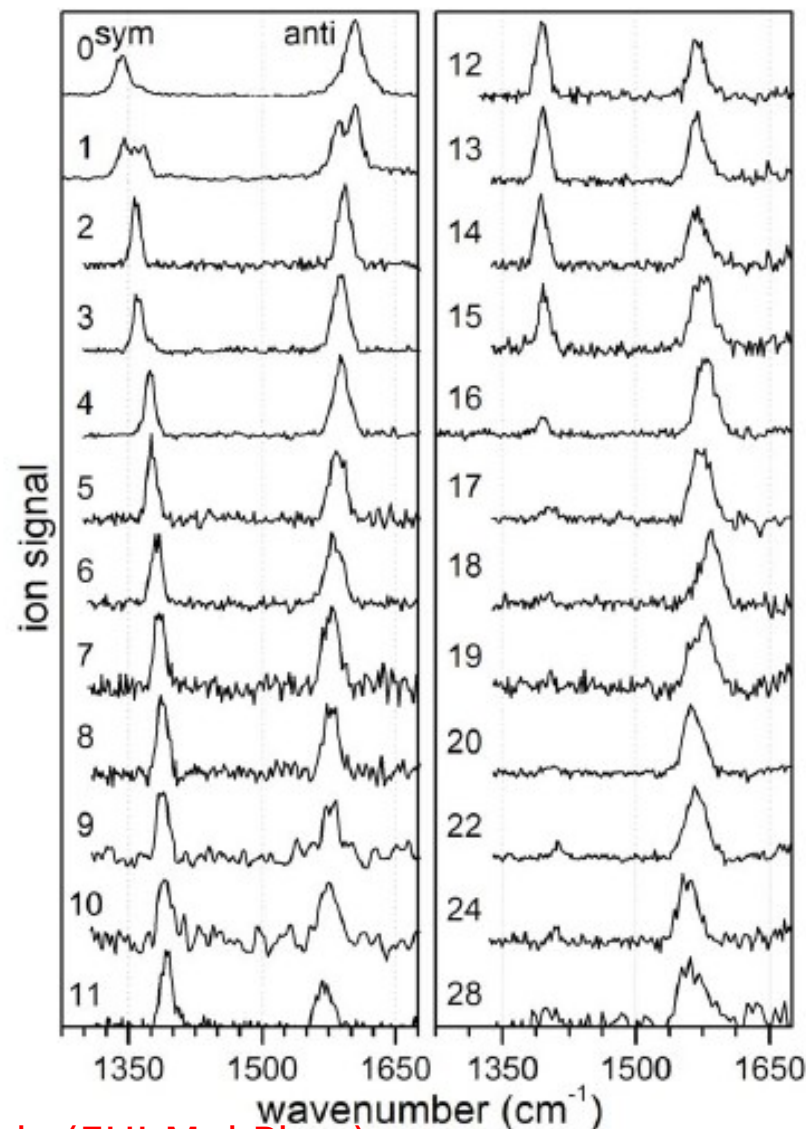
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Microsolvation of the Suberate Dianion



- PE experiments [Yang2004,Wang2008]:
SA folds for $n \geq 14 \dots 16$ water mol.
- IR spectra (PD) at 15 K for clusters of $n=0 \dots 28$ water mol.
- Intensity drop of symmetric C-OO stretching band at $n=16$

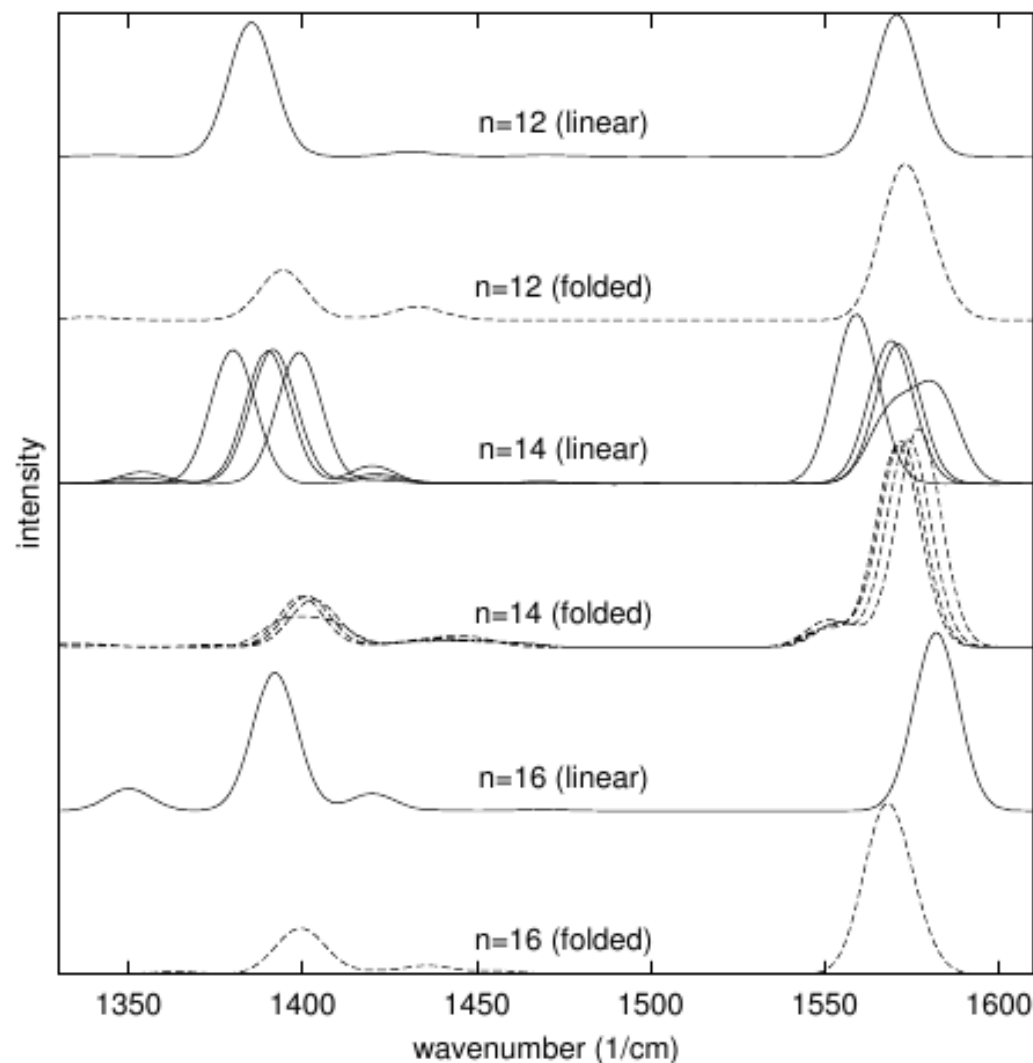
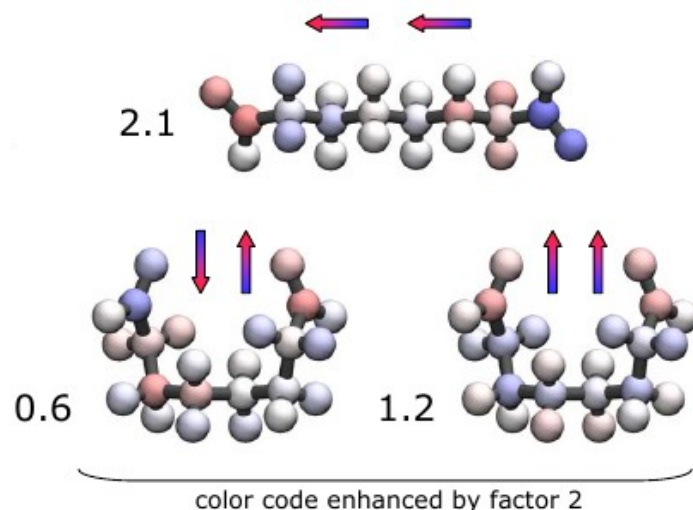


In collaboration with the group of Dr. Knut Asmis (FHI-Mol.Phys)



Symmetric Mode: Intensity Drop upon Folding

- intensity drop due to folding is theoretically reproduced
 - for $n=12...16$
 - for different minimum energy configurations
- ⇒ common feature



Folding Energy

